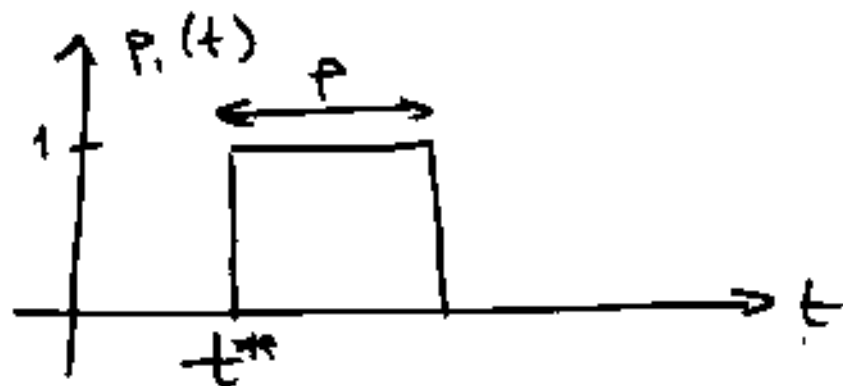


Analysis of Sampled Data:

A pulse can be mathematically described as:



$$P_1(t) = \epsilon(t - t^*) - \epsilon(t - t^* - p)$$

A pulse train starting at time zero can be written as an infinite sum of such pulses:



$$P_2(t) = \sum_{k=0}^{\infty} \{ \epsilon(t - kT) - \epsilon(t - kT - p) \} ; p < T$$

Sometimes, it is preferable to let the pulse train start at time $-p$:

$$P_3(t) = \sum_{k=-\infty}^{+\infty} \left\{ \varepsilon(t - kT) - \varepsilon(t - kT - p) \right\}; \quad p < T$$

$P_3(t)$ is periodic in T , hence we can develop into a Fourier series:

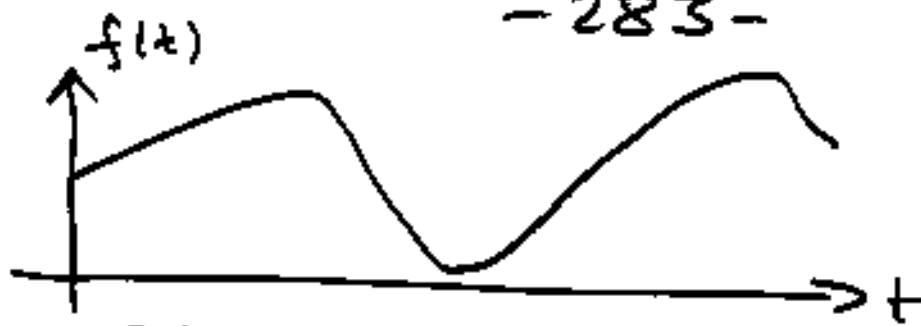
$$P_3(t) = \sum_{n=-\infty}^{+\infty} C_n e^{jn\omega_s t} \quad ; \quad \omega_s = \frac{2\pi}{T}$$

$$C_n = \frac{1}{T} \int_0^T P_3(\tau) \cdot e^{-jn\omega_s \tau} d\tau$$

$$= \frac{1}{T} \int_0^p e^{-jn\omega_s \tau} d\tau = \frac{1 - e^{-jn\omega_s p}}{jn\omega_s T}$$

$$= \left(\frac{p}{T}\right) \cdot \text{sinc}\left(n\omega_s \frac{p}{2}\right) \cdot e^{-jn\omega_s \frac{p}{2}}$$

$$\Rightarrow P_3(t) = \frac{p}{T} \cdot \sum_{n=-\infty}^{+\infty} \text{sinc}\left(n\omega_s \frac{p}{2}\right) \cdot e^{jn\omega_s(t - \frac{p}{2})}$$



$$f_{p_2}^*(t) = f(t) \cdot p_2(t) = \sum_{n=-\infty}^{+\infty} C_n f(t) \cdot e^{jn\omega_s t}$$

$f_{p_2}^*(t)$ is non-periodic. It has a Fourier transform:

$$\begin{aligned} F_{p_2}^*(j\omega) &= \int_{-\infty}^{+\infty} f_{p_2}^*(\tau) e^{-j\omega\tau} d\tau \\ &= \int_{-\infty}^{+\infty} \left(\sum_{n=-\infty}^{+\infty} C_n f(\tau) \cdot e^{jn\omega_s \tau} \right) \cdot e^{-j\omega\tau} d\tau \\ &= \sum_{n=-\infty}^{+\infty} C_n \left\{ \int_{-\infty}^{+\infty} f(\tau) e^{jn\omega_s \tau} \cdot e^{-j\omega\tau} d\tau \right\} \end{aligned}$$

Assuming that $f(t)$ has a Fourier transform, it will be:

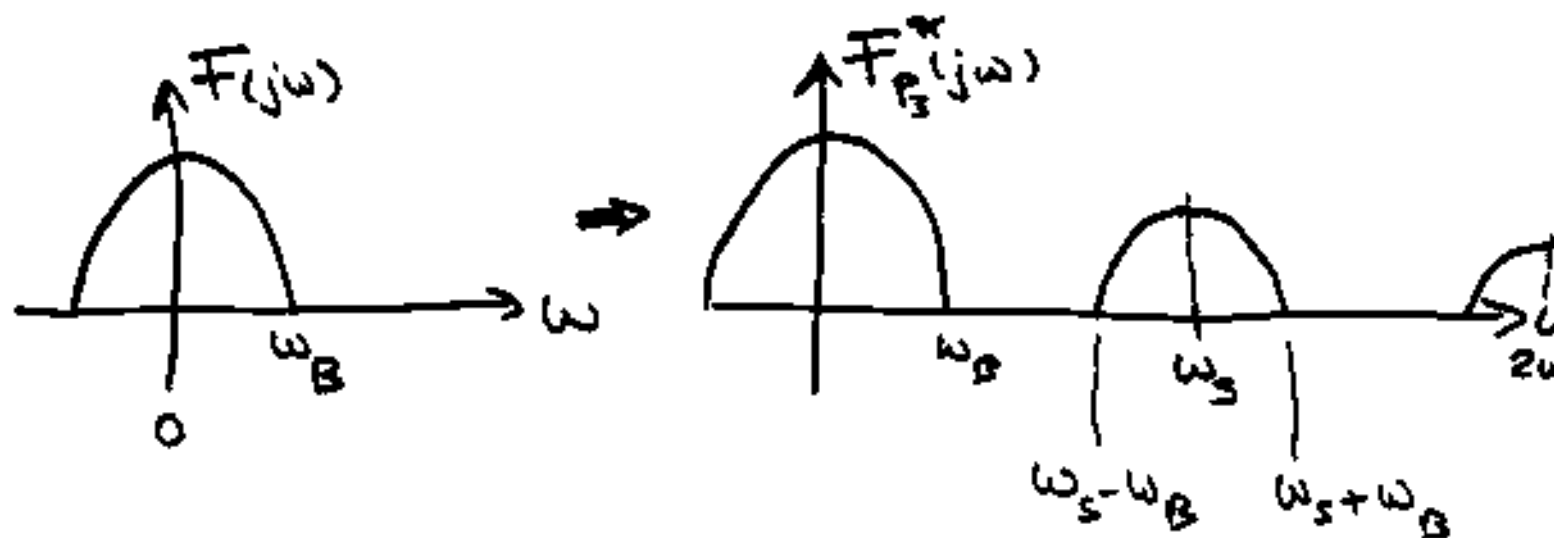
$$F(j\omega) = \mathcal{F}\{f(t)\} = \int_{-\infty}^{+\infty} f(\tau) \cdot e^{-j\omega\tau} d\tau$$

By the shifting theorem:

$$\int_{-\infty}^{+\infty} f(\tau) e^{-j\omega\tau} \cdot e^{jn\omega_s\tau} d\tau \equiv F(j\omega - jn\omega_s)$$

$$\Rightarrow F_p^*(j\omega) = \sum_{n=-\infty}^{+\infty} C_n \cdot F(j\omega - jn\omega_s)$$

Assuming that $F(j\omega)$ is bandlimited, we can draw $F_p^*(j\omega)$ as follows:



⇒ A lowpass filter can retrieve the original signal iff

$$\underline{\underline{\omega_s \geq 2\omega_B}}$$

Shannon's Sampling Theorem.

This provides an upper limit on the sampling rate:

$$\underline{\underline{T = \frac{2\pi}{\omega_s} \leq \frac{\pi}{\omega_B}}}$$

⇒ The sampling rate is now bounded from below and from above:

$$\overbrace{T_{min}} \leq T \leq \overbrace{T_{max}}$$

Dictated by technology in use

Dictated by Systems theory (information/stability)

In practice, it may be more convenient to work with $p_2(t)$ and Laplace transform:

$$F_{p_2}^*(s) = \mathcal{L} \{ f_{p_2}^*(t) \} = \mathcal{L} \{ f(t) \cdot p_2(t) \}$$

$$\Rightarrow F_{p_2}^*(s) = F(s) * P_2(s)$$

↑ convolution

$$p_2(t) = \sum_{k=0}^{\infty} \{ \epsilon(t - kT) - \epsilon(t - kT - p) \}$$

$$\Rightarrow p_2(s) = \sum_{k=0}^{\infty} \left\{ \frac{1}{s} e^{-kTs} - \frac{1}{s} e^{-(kT+p)s} \right\}$$

$$= \sum_{k=0}^{\infty} \frac{1 - e^{-ps}}{s} \cdot e^{-kTs}$$

$$= \frac{1 - e^{-ps}}{s} \underbrace{\left[1 + e^{-Ts} + e^{-2Ts} + \dots \right]}_{H(s)}$$

$$H(s) = 1 + e^{-Ts} + e^{-2Ts} + e^{-3Ts} + \dots$$

$$\Rightarrow e^{-Ts} \cdot H(s) = e^{-Ts} + e^{-2Ts} + e^{-3Ts} + \dots$$

$$\Rightarrow (1 - e^{-Ts}) \cdot H(s) = 1$$

$$\Rightarrow H(s) = \frac{1}{1 - e^{-Ts}}$$

$$\Rightarrow \underline{\underline{P_2(s) = \frac{1 - e^{-ps}}{s(1 - e^{-Ts})}}}}$$

$$f_{P_2}^*(t) = \sum_{k=0}^{\infty} f(kT) \{ \epsilon(t - kT) - \epsilon(t - kT - p) \}$$

$$\approx \sum_{k=0}^{\infty} f(kT) \cdot \{ \epsilon(t - kT) - \epsilon(t - kT - p) \}$$

; $p \ll T$

$$\Rightarrow \underline{\underline{F_{P_2}^*(s) \approx \sum_{k=0}^{\infty} \frac{1 - e^{-ps}}{s} \cdot f(kT) \cdot e^{-kTs}}}}$$

$$= \frac{1 - e^{-ps}}{s} \cdot \sum_{k=0}^{\infty} f(kT) \cdot e^{-kTs}$$

If an ideal sampler would have been used:

$$f^*(t) = \sum_{k=0}^{\infty} f(kT) \cdot \delta(t - kT)$$

$$\Rightarrow F^*(s) = \sum_{k=0}^{\infty} f(kT) \cdot e^{-Ts}$$

$$\Rightarrow F_{p3}^*(s) \approx \frac{1 - e^{-ps}}{s} \cdot F^*(s)$$

If $p \ll T$:

$$1 - e^{-ps} = 1 - \underbrace{\left(1 - ps + \frac{p^2 s^2}{2} - \frac{p^3 s^3}{6} + \dots \right)}_{\text{Taylor series}}$$

$$\approx 1 - 1 + ps = ps$$

$$\Rightarrow F_{p2}^*(s) \approx p \cdot F^*(s)$$

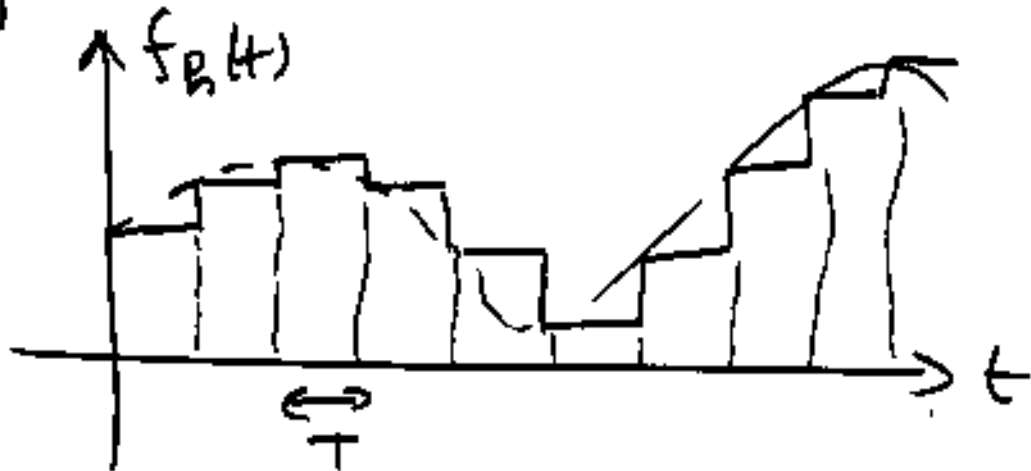
Notice that in the limit
 $p \rightarrow 0$:

$$\lim_{p \rightarrow 0} F_p^*(s) = \phi \neq F^*(s)$$

This is because the amplitude of the pulses was left at 1, rather than $\frac{1}{p}$, as the ideal sampler would.

($\delta(t)$ are pulses of infinitely small width, but infinitely tall amplitude.)

Let us now look at the signal $f_R(t)$:



$$f_R(t) = \sum_{k=0}^{\infty} f(kT) \{ \varepsilon(t - kT) - \varepsilon(t - (k+1)T) \}$$

$$\Rightarrow F_R(s) = \sum_{k=0}^{\infty} f(kT) \cdot \left\{ \frac{1}{s} e^{-kTs} - \frac{1}{s} e^{-(k+1)Ts} \right\}$$

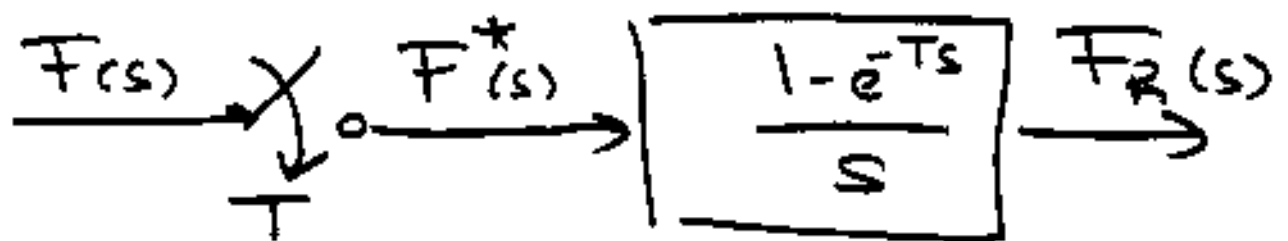
$$= \sum_{k=0}^{\infty} f(kT) \cdot \frac{1}{s} \cdot (1 - e^{-Ts}) \cdot e^{-kTs}$$

$$= \frac{1 - e^{-Ts}}{s} \cdot \underbrace{\sum_{k=0}^{\infty} f(kT) \cdot e^{-kTs}}_{F^*(s)}$$

$F^*(s)$

$$\Rightarrow F_h(s) = \frac{1 - e^{-Ts}}{s} \cdot F^*(s)$$

Graphically:



ideal
Sampler

Laplace transform
of ZOH.

Relationship Between $F(s)$ and $F^*(s)$

Example: $F(s) = \frac{1}{s} \rightarrow f(t) = \epsilon(t)$

$$\Rightarrow f^*(t) = \sum_{k=0}^{\infty} f(kT) \cdot \delta(t - kT)$$

$$= \sum_{k=0}^{\infty} \delta(t - kT)$$

$$\Rightarrow F^*(s) = \sum_{k=0}^{\infty} e^{-kTs} = \frac{1}{1 - e^{-Ts}}$$

Example: $F(s) = \frac{1}{s+a} \Rightarrow f(t) = e^{-at} \cdot \epsilon(t)$

$$\Rightarrow f^*(t) = \sum_{k=0}^{\infty} e^{-akT} \cdot \delta(t - kT)$$

$$\Rightarrow F^*(s) = \sum_{k=0}^{\infty} e^{-kT(s+a)} = \frac{1}{1 - e^{-T(s+a)}}$$

Example: $F(s) = \frac{1}{s(s+1)}$

We use PFE:

$$F(s) = \frac{1}{s} - \frac{1}{s+1}$$

$$\Rightarrow f(t) = \varepsilon(t) - e^{-t} \cdot \varepsilon(t)$$

$$\begin{aligned} \Rightarrow F^*(s) &= \frac{1}{1 - e^{-Ts}} - \frac{1}{1 - e^{-T(s+1)}} \\ &= \frac{(1 - e^{-T} \cdot e^{-Ts}) - (1 - e^{-Ts})}{(1 - e^{-Ts}) \cdot (1 - e^{-T} \cdot e^{-Ts})} \\ &= \frac{(1 - e^{-T}) \cdot e^{-Ts}}{(1 - e^{-Ts}) \cdot (1 - e^{-T} \cdot e^{-Ts})} \end{aligned}$$

Physical interpretation:

$S \hat{=} \text{derivative } \left(\frac{d}{dt} \right)$

$S^{-1} = S^{-1} \hat{=} \text{integral } \left(\int \right)$

$Z \hat{=} \text{left shift by } T$

$\frac{1}{Z} = Z^{-1} \hat{=} \text{right shift by } T$

